

Mixed normal-superconducting states in the presence of strong electric currents

Yaniv Almog, Bernard Helffer, Xingbin Pan

LSU, Orsay

Henri Lebesgue Center

Superconductivity

Onnes (1911), Meissner & Ochsenfeld (1933)

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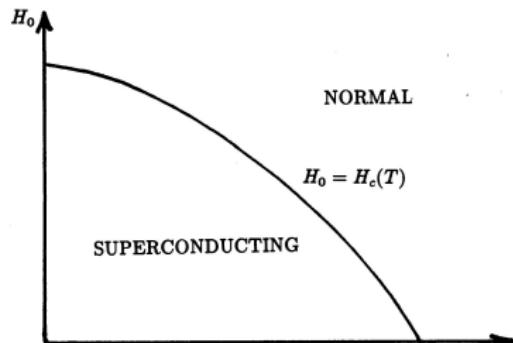


FIG. 1. *Critical field as a function of temperature.*

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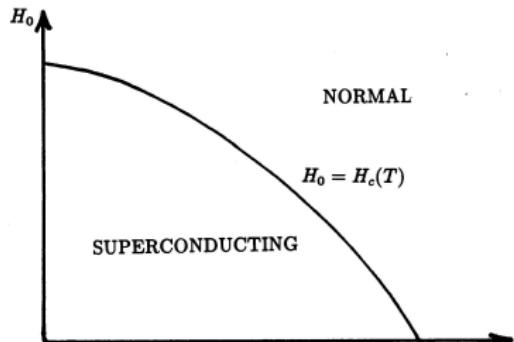


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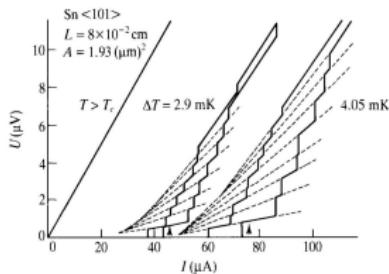


FIGURE 11.11
Current-voltage characteristics of tin "whiskers" showing regular step structures due to successive establishment of phase-slip centers. Here $\Delta T = T_c - T$. (After Meyer.)





$$\frac{\partial \psi_\kappa}{\partial t} - \nabla_{\kappa A_\kappa}^2 \psi_\kappa + i\kappa \phi_\kappa \psi_\kappa = \kappa^2 (1 - |\psi_\kappa|^2) \psi_\kappa \quad \text{in } (0, +\infty) \times \Omega,$$

$$\frac{\partial A_\kappa}{\partial t} + \nabla \phi_\kappa + \operatorname{curl}^2 A_\kappa = \frac{1}{\kappa} \Im(\bar{\psi}_\kappa \nabla_{\kappa A_\kappa} \psi_\kappa) \quad \text{in } (0, +\infty) \times \Omega,$$

$$\psi = 0 \quad \text{on } (0, +\infty) \times \partial\Omega_c,$$

$$\nabla_{\kappa A_\kappa} \psi_\kappa \cdot \nu = 0 \quad \text{on } (0, +\infty) \times \partial\Omega_i,$$

$$\frac{\partial \phi_\kappa}{\partial \nu} = -\kappa J(x) \quad \text{on } (0, +\infty) \times \partial\Omega,$$

$$\int_{\partial\Omega} \operatorname{curl} A_\kappa \, ds = \kappa h_{\text{ex}} \quad \text{on } (0, +\infty),$$

Gauge invariance

$$A \rightarrow A + \nabla\omega \quad \phi \rightarrow \phi - \frac{\partial\omega}{\partial t} \quad \psi \rightarrow \psi e^{i\omega}$$

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$$J|_{\partial\Omega_i} = 0 \Rightarrow B_\kappa|_{\partial\Omega_{i,j}} \stackrel{\text{def}}{=} \operatorname{curl} A_\kappa|_{\partial\Omega_{i,j}} = h_j \kappa \quad J|_{\partial\Omega_c} \neq 0 \quad j \in \{1, 2\}$$

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$$h_1 h_2 < 0,$$

$$h = \max(|h_1|, |h_2|)$$

$$1 < h \leq \frac{1}{\Theta_0} \quad \text{or} \quad \frac{1}{\Theta_0} < h \ll \kappa$$

$$\Theta_0 \sim 0.59 \text{ de Gennes' value}$$

Linearized equation

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Normal fields: $(\psi, A, \phi) = (0, \kappa A_n, c\kappa\phi_n)$ $c = \kappa^2/\sigma$

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$$\mathcal{L}_\kappa = (i\nabla + \kappa^2 A_n)^2 + i c \kappa^2 \phi_n$$

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$$\mu(\kappa, J) \stackrel{\text{def}}{=} \inf_{\lambda \in \sigma(\mathcal{L}_\kappa)} \Re \lambda < 1 \Rightarrow u \equiv 0 \text{ unstable}$$

$$J = J_0 \kappa \Rightarrow h \sim \mathcal{O}(\kappa)$$

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$$\mathcal{A}_+(z_i) = -\left(\nabla - ij_i \frac{s^2}{2} \hat{i}_t\right)^2 + icj_i t$$

$$D(\mathcal{A}^+) = \{u \in H_0^{1,\text{mag}}(\mathbb{R}_+^2, \mathbb{C}) \cap L^2(\mathbb{R}_+^2, \mathbb{C}; y \, dx dy) : \mathcal{A}_+ u \in L^2(\mathbb{R}_+^2, \mathbb{C})\}$$

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Theorem

Let $\mu_\infty = \liminf_{\kappa \rightarrow \infty} \mu$. Then

$$\limsup_{\kappa \rightarrow \infty} \mu = \mu_\infty = \nu_m.$$

Almog, Helffer (2014)

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(1) *satisfied* $\forall K \subset \omega_j \cup \partial\Omega_{i,j} \cup \partial\Omega_c$ *compact*



$$\|\nabla_{\kappa A_\kappa}(\eta\psi_\kappa)\|_2^2 \leq \kappa^2 \|\eta\psi_\kappa\|_2^2 - \|\psi_\kappa \nabla \eta\|_2^2$$

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$$K \cap \partial\Omega_i \neq \emptyset \Rightarrow \kappa^2 \Theta_{\pi/2}^{DN} \inf_{x \in K} |B_n(x)| \|\eta\psi_\kappa\|_2^2 \leq \|\nabla_{\kappa^2 A_n}(\eta\psi_\kappa)\|_2^2$$

Dirichlet-Neumann corner: Dauge

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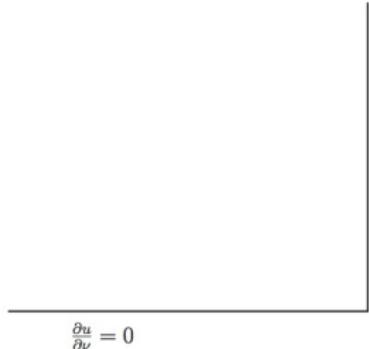
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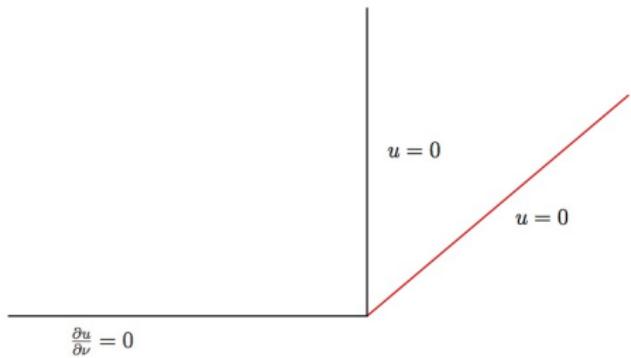
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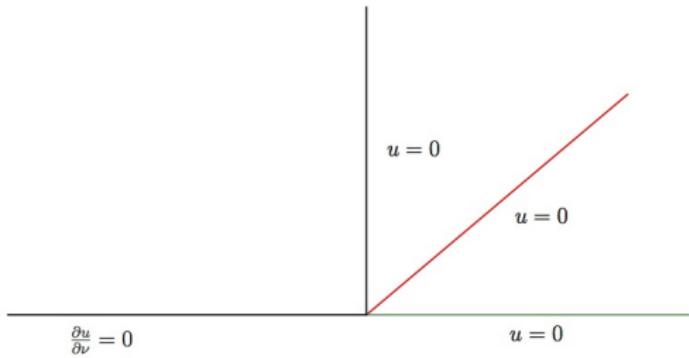


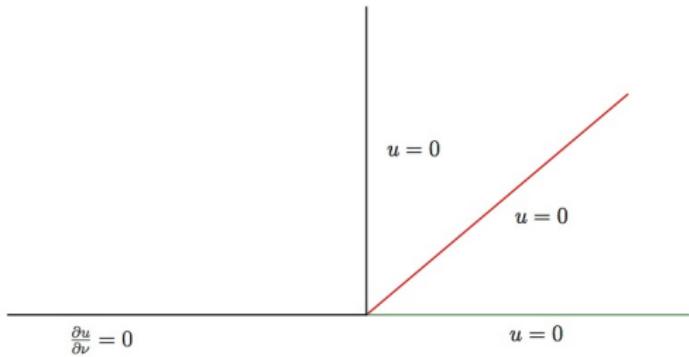


$u = 0$

$$\frac{\partial u}{\partial \nu} = 0$$







$$\Theta_0 \geq \Theta_{\alpha}^{DN} \geq \Theta_{\pi}^{DN} \geq \Theta_0$$

Proposition

$$\exists C(J, \Omega) > 0 : \forall \kappa \geq 1 \|\psi_\kappa\|_2 \leq C(J, \Omega) (1 + c^{-1/2})^{1/3} \kappa^{-1/6}$$

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$$E_{min} \sim \kappa^2 \left[\frac{1}{2} + \int_{\Omega} g(B_n(x)) dx \right]$$

$-1/2 < g(b) \leq g(1) = 0$ increasing - Sandier & Serfaty (2003)

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Theorem

For any compact set $K \subset \omega_j \cup \partial\Omega_c$ there exist $C > 0$, $\alpha > 0$, and $\kappa_0 \geq 1$, such that for any $\kappa \geq \kappa_0$ we have

$$\limsup_{t \rightarrow \infty} \int_K |\psi_\kappa(t, x)|^2 dx \leq Ce^{-\alpha\kappa} .$$

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$$\begin{aligned} \Delta_A \psi + \psi (1 - |\psi|^2) - i\phi\psi &= 0 && \text{in } \Omega^R, \\ \operatorname{curl}^2 A + \nabla\phi &= \Im(\bar{\psi} \nabla_A \psi) && \text{in } \Omega^R, \\ \psi &= 0 && \text{on } \partial\Omega_c^R, \\ \nabla_A \psi \cdot \nu &= 0 && \text{on } \partial\Omega_i^R, \\ \frac{\partial\phi}{\partial\nu} &= R^{-(1-\gamma)} J && \text{on } \partial\Omega^R, \\ \int_{\partial\Omega^R} \operatorname{curl} A(x) ds &= R^\gamma h_{ex}. \end{aligned}$$

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$$0 < \gamma < 1$$

Proposition

$\exists K \subset \Omega$ *compact*, $C > 0$, $R_0 > 0$, $\alpha > 0$:

$$R > R_0 \Rightarrow \int_{K_R} |\psi(x)|^2 dx \leq Ce^{-\alpha R},$$

$$K_R = T_R(K)$$

Proposition

$\exists K \subset \Omega$ *compact*, $C > 0$, $R_0 > 0$, $\alpha > 0$:

$$R > R_0 \Rightarrow \int_{K_R} |\psi(x)|^2 dx \leq Ce^{-\alpha R},$$

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$$\Delta_{\epsilon^{-2}A_\epsilon} \psi_\epsilon + \frac{\psi_\epsilon}{\epsilon^2} (1 - |\psi_\epsilon|^2) - \frac{i}{\epsilon^2} \phi_\epsilon \psi_\epsilon = 0 \quad \text{in } \Omega$$

$$\operatorname{curl}^2 A_\epsilon + \nabla \phi_\epsilon = \Im(\bar{\psi}_\epsilon \nabla_{\epsilon^{-2}A_\epsilon} \psi_\epsilon) \quad \text{in } \Omega$$

$$\frac{\partial \phi_\epsilon}{\partial \nu} = \epsilon^{-\gamma} J \quad \text{on } \partial \Omega$$

$$A_{1,\epsilon} = A_\epsilon - \epsilon^{-\gamma} A_n \quad ; \quad B_{1,\epsilon} = \operatorname{curl} A_{1,\epsilon} .$$

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Conjecture

For any compact set $K \subset \Omega \setminus B_n^{-1}(0)$, $\exists R_0 > 0$, $C > 0$, and $\alpha > 0$

$$R > R_0 \Rightarrow \int_{K_R} |\psi(x)|^2 dx \leq C e^{-\alpha R}$$

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Table 9.6.

Critical current density J_c for various superconductors. Values for the temperature T_{max} and the field B_{app} of measurement are also given. References are identified by first author and year.

Material	J_c (A/cm ²)	B_{app} (tesla)	T_{max} (K)	Comments	Reference
Melt textured high- T_c materials		0.02 to 1	77	$J_c \sim [1 - (T/T_c)^\alpha]/B^\beta$, where $\alpha \leq 7$, $\beta \sim 0.5$	Fisher (1994)
Cu-Nb composite	2.5×10^5	Self field	4.2	$T_c = 8.8$ K	Reddy (1986)
NbTi	3.4×10^5	5	4.2	Superconducting composite	Li (1983)
	3×10^5	5		Filament diameters 5 to 9 μm	Kanithi (1989)
	3×10^5	0.5	4.2	Filament diameter 50–500 nm; see $\zeta\lambda$ table	Cave (1989)
	1.41×10^6	1		Wire enclosed by CuNi alloy, 0.072 μm filament diameter, J_c is less for 0.061 μm filament diameter	Kumano (1990)
	1.07×10^6	1		0.4 μm filament, artificial pinning control	Miura (1992)
	4.44×10^5	3		0.4 μm filament, artificial pinning control	Miura (1992)
	5.2×10^5	5	4.2	47 wt% Ti, filaments are aligned	Cooley (1991)
	2.85×10^5	5		48 wt% Ti, 210 filament, 1.6 mm diameter wire	Cherny (1997)
	3.5×10^5	5	4.2	Industrial wires	Chernoplekov (1992)
	3.36×10^5	3		Wire, 54 wt% Ti	Inoue (1995)
	4.6×10^5	5	4.2	Artificial pinning center wire, 47 wt% Ti	Heussner (1997)
NbN	2×10^4	0		Josephson junction 1.5 μm^2 area	Tarutani (1984)

(continued)